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# Erratum: Quantum propagation of neutral atoms in a magnetic quadrupole guide [Phys. Rev. A 61, 033614 (2000)]

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The quantum number  $\ell$  in the wave function  $\psi(\rho, \phi)$  is in fact not an eigenvalue of the orbital angular momentum  $L_z$  of the motion of the atom around the center of the guide, as we had mistakenly assumed without checking, but an eigenvalue of the operator  $L_z - s_z$ , where  $s_z$  is the  $z$  component of the total internal angular momentum of the atom that couples to the magnetic field. For a magnetic quadrupole field of the form of Eq. (1.1),  $\mathbf{B} = (4B_0x/R, -4B_0y/R, 0)$  the operators  $L_z$  and  $s_z$  do not individually commute with the Hamiltonian  $H = \mathbf{p}^2/(2m) + g\mu_B \mathbf{s} \cdot \mathbf{B}$ ,

$$[L_z, H] = ig\mu_B \hbar \frac{4B_0}{R} (s_y x + s_x y),$$

$$[s_z, H] = ig\mu_B \hbar \frac{4B_0}{R} (s_y x + s_x y),$$

but the difference  $L_z - s_z$  commutes with  $H$  and is thus a conserved quantity,  $[L_z - s_z, H] = 0$ . Calculation shows that the quantum number  $\ell$  in the wave function  $\psi(\rho, \phi)$  is the eigenvalue of  $L_z - s_z$ , e.g., for spin 1/2

$$(L_z - s_z) \frac{1}{\sqrt{2}} \begin{pmatrix} F_+(\rho) e^{i(\ell+1/2)\phi} \\ F_-(\rho) e^{i(\ell-1/2)\phi} \end{pmatrix} = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} (\ell+1/2) F_+(\rho) e^{i(\ell+1/2)\phi} \\ (\ell-1/2) F_-(\rho) e^{i(\ell-1/2)\phi} \end{pmatrix} - \frac{\hbar}{2\sqrt{2}} \begin{pmatrix} F_+(\rho) e^{i(\ell+1/2)\phi} \\ -F_-(\rho) e^{i(\ell-1/2)\phi} \end{pmatrix}$$

$$= \ell \frac{\hbar}{\sqrt{2}} \begin{pmatrix} F_+(\rho) e^{i(\ell+1/2)\phi} \\ F_-(\rho) e^{i(\ell-1/2)\phi} \end{pmatrix},$$

and similarly for the spin 1 wave function in Eq. (4.5). Although  $L_z$  and  $s_z$  do not commute with  $H$ , they commute with  $L_z - s_z$  and with each other, which means that any eigenstate of  $L_z - s_z$  must be a linear combination of simultaneous eigenstates of  $L_z$  and  $s_z$ . Since the eigenvalues of  $L_z$  are always integer, it follows that the eigenvalue  $\ell$  of  $L_z - s_z$  must be half integer for half integer spin  $s$  and integer for integer spin  $s$ . Only then does the wave function have the correct transformation properties under rotations. For spin 1/2 a rotation by  $2\pi$  around the  $z$  axis, for example, transforms the wave function of the above eigenstate into

$$\frac{1}{\sqrt{2}} F_+(\rho) e^{i(\ell+1/2)(\phi+2\pi)} \mathcal{D}_{mm'}^{1/2}(0,0,2\pi) \begin{pmatrix} 1 \\ 0 \end{pmatrix} + \frac{1}{\sqrt{2}} F_-(\rho) e^{i(\ell-1/2)(\phi+2\pi)} \mathcal{D}_{mm'}^{1/2}(0,0,2\pi) \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

Since the application of the spin rotation matrix  $\mathcal{D}_{mm'}^{1/2}(0,0,2\pi)$  causes the spinors to change sign, the  $\phi$  dependent prefactors must not change sign under  $\phi \rightarrow \phi + 2\pi$ . Thus  $\ell$  has to be half integer for spin 1/2, and not integer as was erroneously stated in Sec. III. It follows that  $\ell$  cannot be zero, which, as explained by Eq. (5.3), has the consequence that there are no exact bound states in the case of spin 1/2 and makes Sec. III A redundant. Solving Eqs. (3.11) for  $\ell = 1/2, 3/2, \dots$  gives results that are qualitatively the same as those of Sec. III B. For  $\ell = 1/2$  the energies and widths ( $\epsilon_i, \Gamma_i$ ) of the first three resonances are (2.64, 0.34), (4.25, 0.34), and (5.62, 0.34), and for  $\ell = 3/2$  they are (3.53, 0.11), (5.04, 0.15), and (6.34, 0.18), in units of  $(\hbar^2 G^2/2m)^{1/3}$ . The similarity of the motion of spin 1/2 and of spin 1 atoms in the guide is underlined by the fact that Eqs. (3.7) for spin 1/2 and  $\ell = 1/2$  are structurally the same as Eqs. (4.7a) and (4.7b) for spin 1 and  $\ell = 0$  except for a replacement of  $G$  by  $2G$ .

The energies and widths for  $s = 1/2$  and the unphysical case of integer  $\ell$  have been reproduced by a complex scaling calculation by Potvliege and Zehnlé [1]. These authors agree that  $\ell$  should in fact be taken half integer [2]. The applicability of their method is independent of whether integer or half integer values of  $\ell$  are being considered.

Nothing changes for spin 1 in Sec. IV.

There are also two misprints. Contrary to what is stated in the text of Sec. II, we have used  $G = 2g\mu_B B_0/R$  throughout the paper. The last paragraph of Sec. VII discusses the first excitation energy for an  $\ell = 0$  atom of  $^{87}\text{Rb}$ , and not  $\ell = 1$  as originally stated.

We are indebted to Dr. John Stockton, Dr. Clifford Hicks, and Professor Hideo Mabuchi for pointing out to us that  $\ell$  is an eigenvalue of  $L_z - s_z$  and must be half integer for spin 1/2.

[1] R.M. Potvliege and V. Zehnlé, Phys. Rev. A **63**, 025601 (2001).

[2] R. M. Potvliege, private communication.